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FREE VIBRATION ANALYSIS OF RECTANGULAR PLATE BY SPLIT-DEFLECTION METHOD

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Abstract

This paper presents free vibration analysis of rectangular plate by split-deflection method. In this method, the deflection was split into x and y components of deflection. That is the deflection of the rectangular plate was taken as the product of these two components. Having made this assumption, the study went ahead to formulate total potential energy functional from principles of theory of elasticity based on work-error approach. This energy functional was minimized by direct variation and equation for resonating frequency was obtained. Two illustrative examples were used to test this method. They are plates with all edges simply supported and all edges clamped. The first case used polynomial function for x component of deflection and trigonometric function for y component of. However, the second example used polynomial function for both x and y components of deflection. Fundamental resonating frequencies (in non dimensional forms) of the two plates for aspect ratios ranging from 1.0 to 2.0 (at increment of 0.1) were determined and compared with the values from previous study. From the comparison, it was observed that the maximum percentage difference of 0.06 was recorded for the first example at aspect ratios of 1.1 and 1.3 with non dimensional resonating frequencies of 18.03 and 15.71 respectively. For the

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second example, the maximum percentage difference of 0.09 was recorded at aspect ratio of 1.1 with non dimensional resonating frequencies of 32.96. This small value of percentage differences show that this present method is adequate and reliable for classical plate theory (CPT) free vibration analysis of rectangular plates.

Index Terms—Resonating frequency, split-deflection, work-error, energy functional, polynomial function, trigonometric function

I. INTRODUCTION

Most energy methods used for classical plate theory (CPT) free vibration analysis of rectangular plates include Raleigh, Raleigh-Rit, Ritz, Galerkin, minimum potential energy, work-erroretc (Ugural, 1999, Ventsel and Krauthammer, 2001 and Ibearugbulem et al., 2014). The deflection (displacement normal to the plane of the plate is a single orthogonal function, w.This is apparent in the energy functional for Raleigh, Raleigh-Ritz and Ritz. Typical Ritz energy functional is (Ibearugbulem et al., 2014):

$$\Pi = \frac{D}{2} \int_0^a \int_0^b \left(\left[\frac{\partial^2 w}{\partial x^2} \right]^2 + 2 \left[\frac{\partial^2 w}{\partial x \partial y} \right]^2 + \left[\frac{\partial^2 w}{\partial y^2} \right]^2 \right) \partial x \partial y$$
$$- \frac{m \cdot \lambda^2}{2} \int_0^a \int_0^b w^2 \, \partial x \partial y$$

The use of single orthogonal deflection function is also seen in Galerkinand work-error methods. Typical work-error functional is (Ibearugbulem et al., 2014):

$$\Pi = \frac{D}{2} \int_0^a \int_0^b \left(\frac{\partial^4 w}{\partial x^4} w + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} w + \frac{\partial^4 w}{\partial y^4} w \right) \partial x \partial y$$
$$- \frac{m \cdot \lambda^2}{2} \int_0^a \int_0^b w^2 \, \partial x \partial y$$

From the literature so far, most scholarly works on CPT analysis of rectangular plates rely on this single orthogonal function (Hutchinson, 1992, Jianqiao, 1994, Ugural, 1999, Ventsel and Krauthammer, 2001, Wang et al., 2002, Taylor and Govindjee, 2004, Szilard, 2004, Jiu et al., 2007, Erdem et al., 2007, Ezeh et al., 2013, Ibearugbulem, 2014). Obviously, one can assert that all energy functional in use are based on single orthogonal deflection function and none has used a

deflections function that is typically separated into two independent distinct functions ($w = w_x * w_y$). In this present studyw_xand w_ymay be both polynomials or trigonometric functions or w_x may be polynomialswhilew_y may be trigonometry. The main reason for this modification is to help the analysis who may have difficulty in obtaining orthogonal function for a plate of a particular boundary condition. In this case, the analyst who may have easy access to deflection equations for beams of any boundary condition can find the proposed method quite useful and handy.

II. BASIC ASSUMPTIONS

The assumption here is that the general deflection, w is split into w_x and w_y . That is the split-deflection function is given as:

$$w = w_{x}.w_{y}$$
 1

Where the w_x and w_y components of the deflection are defined as:

$$w_x = \sqrt{A}.\,h_1$$

$$w_y = \sqrt{A}.h_2$$

Substituting equations (2) and (3) into equation (1) gives:

$$w = A h_1 h_2 4$$

III. IN-PLANE DISPLACEMENTS

From the assumption that vertical shear strains are zero for classical plate and making use of split-deflection, we obtain:

$$u = -z\frac{dw}{dx} = -z\frac{dw_x}{dx}w_y 5$$

$$v = -z \frac{dw}{dy} = -z \frac{dw_y}{dy} w_x \tag{6}$$

IV. STRAIN DEFLECTION RELATIONSHIP

Differentiating equations (5) and (6), the three in-plane strains of CPT are given as:

$$\varepsilon_x = \frac{du}{dx} = -z \frac{d^2 w_x}{dx^2} w_y \tag{7}$$

$$\varepsilon_{y} = \frac{dv}{dy} = -z \frac{d^{2}w_{y}}{dy^{2}} w_{x}$$
 8

$$\gamma_{xy} = \frac{du}{dy} + \frac{dv}{dx} = -2z \frac{dw_x}{dx} \frac{dw_y}{dy}$$

V. STRESS-STRAIN RELATIONSHIP

The CPT constitutive equations for plane stress plate are given as:

$$\sigma_{x} = \frac{E}{1 - \mu^{2}} \left[\varepsilon_{x} + \mu \varepsilon_{y} \right]$$
 10

$$\sigma_{y} = \frac{E}{1 - \mu^{2}} \left[\mu \varepsilon_{x} + \varepsilon_{y} \right]$$
 11

$$\tau_{xy} = \frac{E(1-\mu)}{2(1-\mu^2)} \gamma_{xy}$$
 12

VI. STRESS - DEFLECTION RELATIONSHIP

Substituting equations (7), (8) and (9) into equations (10), (11) and (12) where appropriate gives the split-deflection stress-deflection equation as:

$$\sigma_{x} = \frac{-Ez}{1 - \mu^{2}} \left[\frac{d^{2}w_{x}}{dx^{2}} w_{y} + \mu \frac{d^{2}w_{y}}{dy^{2}} w_{x} \right]$$
 13

$$\sigma_{y} = \frac{-Ez}{1 - \mu^{2}} \left[\mu \frac{d^{2}w_{x}}{dx^{2}} w_{y} + \frac{d^{2}w_{y}}{dy^{2}} w_{x} \right]$$
 14

$$\tau_{xy} = \frac{-Ez(1-\mu)}{(1-\mu^2)} \frac{dw_x}{dx} \frac{dw_y}{dy}$$
 15

VII. TOTAL POTENTIAL ENERGY

The strain energy is defined as:

$$U = \frac{1}{2} \int_{x} \int_{y} \left[\int_{-\frac{t}{2}}^{\frac{t}{2}} [\sigma_{x} \varepsilon_{x} + \sigma_{x} \varepsilon_{x} + \sigma_{x} \varepsilon_{x}] dz \right] dx dy \quad 16$$

For pure bending analysis, the external work is given as:

$$V = \int_{x} \int_{y} \frac{m \cdot \lambda^{2}}{2} w_{x}^{2} \cdot w_{y}^{2} dx dy$$

That is

$$V = \frac{m \cdot \lambda^{2}}{2} \int_{x} w_{x}^{2} dx \int_{y} w_{y}^{2} dy$$
 17

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Substituting equations (10) to (15) into equation (16) gives strain energy – deflection relationship as:

$$U = \frac{D}{2} \int_{x} \int_{y} \left[\left(\frac{d^{2} w_{x}}{dx^{2}} \right)^{2} w_{y}^{2} + 2 \left(\frac{dw_{x}}{dx} \right)^{2} \left(\frac{dw_{y}}{dy} \right)^{2} + \left(\frac{d^{2} w_{y}}{dy^{2}} \right)^{2} w_{x}^{2} \right] dx dy$$

That is in the work-error approach, the strain energy becomes:

$$U = \frac{D}{2} \left[\int_{x} \frac{d^{4}w_{x}}{dx^{4}} w_{x} dx \int_{y} w_{y}^{2} dy \right]$$

$$+ \frac{2D}{2} \left[\int_{x} \frac{d^{2}w_{x}}{dx^{2}} w_{x} dx \int_{y} \frac{d^{2}w_{y}}{dy^{2}} w_{y} dy \right]$$

$$+ \frac{D}{2} \left[\int_{x} w_{x}^{2} dx \int_{y} \frac{d^{4}w_{y}}{dy^{4}} w_{y} dy \right]$$
18

Subtracting equation (17) from Equation (18) gives the total potential energy functional as:

$$\Pi = \frac{D}{2} \left[\int_{x} \frac{d^{4}w_{x}}{dx^{4}} w_{x} dx \int_{y} w_{y}^{2} dy \right]
+ \frac{2D}{2} \left[\int_{x} \frac{d^{2}w_{x}}{dx^{2}} w_{x} dx \int_{y} \frac{d^{2}w_{y}}{dy^{2}} w_{y} dy \right]
+ \frac{D}{2} \left[\int_{x} w_{x}^{2} dx \int_{y} \frac{d^{4}w_{y}}{dy^{4}} w_{y} dy \right] - \frac{m \cdot \lambda^{2}}{2} \int_{x} w_{x}^{2} dx \int_{y} w_{y}^{2} dy$$
19

Substituting equations (1) and (2) into equation (19) gives:

$$\Pi = \frac{A^2 D}{2} \left[\int_x \frac{d^4 h_1}{dx^4} h_1 \, dx \int_y h_2^2 \, dy \right]
+ \frac{2A^2 D}{2} \left[\int_x \frac{d^2 h_1}{dx^2} h_1 \, dx \int_y \frac{d^2 h_2}{dy^2} h_2 \, dy \right]
+ \frac{A^2 D}{2} \left[\int_x h_1^2 \, dx \int_y \frac{d^4 h_2}{dy^4} h_2 \, dy \right]
- \frac{m \cdot \lambda^2}{2} A^2 \int_x h_1^2 dx \int_y h_2^2 \, dy$$
20

Now, equation (20) can be written in non dimensional axes R and Q.

$$x = aR$$

$$y = aQ$$

$$P = b/a$$
21
22

Where a, b and P are the plate lengths in x and y axes and long span- short span aspect ratio respectively.

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Substituting equations (21), (22) and (23) into equation (20) gives:

$$\begin{split} \Pi &= \frac{abA^2D}{2a^4} \left[\int_0^1 \frac{d^4h_1}{dR^4} \, h_1 dR \int_0^1 h_2^2 \, dQ \right] \\ &+ 2 \frac{abA^2D}{2a^4P^2} \left[\int_0^1 \frac{d^2h_1}{dR^2} \, h_1 dR \int_0^1 \frac{d^2h_2}{dQ^2} \, h_2 \, dQ \right] + \frac{abA^2D}{2a^4P^4} \left[\int_0^1 h_1^2 \, dR \int_0^1 \frac{d^4h_2}{dQ^4} \, h_2 \, dQ \right] \\ &- \frac{m \cdot \lambda^2}{2} A^2 ab \int_0^1 h_1^2 dR \int_0^1 h_2^2 \, dQ \end{split} \qquad \qquad 24 \end{split}$$

VIII. DIRECT VARIATION OF TOTAL POTENTIAL ENERGY

Equation (24) shall be differentiated with respect to the deflection coefficient, A and the result is:

$$\begin{split} \frac{d\Pi}{dA} &= \frac{AD}{a^4} \left[\int_0^1 \frac{d^4 h_1}{dR^4} h_1 dR \int_0^1 h_2^2 dQ \right] \\ + 2 \frac{AD}{a^4 P^2} \left[\int_0^1 \frac{d^2 h_1}{dR^2} h_1 dR \int_0^1 \frac{d^2 h_2}{dQ^2} h_2 dQ \right] \\ + \frac{AD}{a^4 P^4} \left[\int_0^1 h_1^2 dR \int_0^1 \frac{d^4 h_2}{dQ^4} h_2 dQ \right] \\ - m \cdot \lambda^2 A \int_0^1 h_1^2 dR \int_0^1 h_2^2 dQ = 0 \end{split}$$

That is

$$\frac{D}{a^4} \left[\int_0^1 \frac{d^4 h_1}{dR^4} h_1 dR \int_0^1 h_2^2 dQ \right]
+ 2 \frac{D}{a^4 P^2} \left[\int_0^1 \frac{d^2 h_1}{dR^2} h_1 dR \int_0^1 \frac{d^2 h_2}{dQ^2} h_2 dQ \right]
+ \frac{D}{a^4 P^4} \left[\int_0^1 h_1^2 dR \int_0^1 \frac{d^4 h_2}{dQ^4} h_2 dQ \right] = m \cdot \lambda^2 \int_0^1 h_1^2 dR \int_0^1 h_2^2 dQ$$
25

This equation (25) is the direct governing equation of rectangular plate under free vibration using work-error approach from this present method. Rearranging equation (25) and making resonating frequency, λ the subject of the equation gives:

$$\lambda^2 = \left(\frac{k_x + 2\frac{k_{xy}}{P^2} + \frac{k_y}{P^4}}{k_\lambda}\right) * \frac{D}{ma^4}$$

Where

$$k_x = \int_0^1 \frac{d^4 h_1}{dR^4} h_1 dR \int_0^1 h_2^2 dQ$$
 27

$$k_{xy} = \int_0^1 \frac{d^2 h_1}{dR^2} h_1 dR \int_0^1 \frac{d^2 h_2}{dQ^2} h_2 dQ$$
 28

$$k_y = \int_0^1 h_1^2 dR \int_0^1 \frac{d^4 h_2}{dQ^4} h_2 dQ$$
 29

$$k_{\lambda} = \int_{0}^{1} h_{1}^{2} dR \int_{0}^{1} h_{2}^{2} dQ$$
 30

IX. NUMERICAL EXAMPLE

Analyze a classical rectangular thin isotropic plate with:

i all the four edge simply supported using polynomial and trigonometry functions respectively for w_x and w_y .

iiall the four edge clamped using only polynomial function for both w_x and w_y

All Simple supported edge rectangular plate $w_x = \sqrt{A} \left(R - 2R^3 + R^4 \right) \qquad \qquad 31$ $w_Y = \sqrt{A} \sin \pi Q \qquad \qquad 32$

From equations (31) and (32), h_1 and h_2 are:

$$h_1 = R - 2R^3 + R^4 33$$

$$h_2 = \sin \pi Q$$
 34

With these we obtain:

$$\int_0^1 h_1^2 dR = \frac{31}{630} \text{ and } \int_0^1 h_2^2 dQ = 0.5.$$

$$\int_0^1 \frac{d^4 h_1}{dR^4} h_1 dR = 4.8 \text{ and } \int_0^1 \frac{d^4 h_2}{dQ^4} h_2 dQ = 0.5 \pi^4$$

$$\int_0^1 \frac{d^2 h_1}{dR^2} h_1 dR = \frac{17}{35} \text{ and } \int_0^1 \frac{d^2 h_2}{dQ^2} h_2 dQ = 0.5 \pi^2.$$

$$k_r = (4.8)(0.5) = 2.4$$
 35

$$k_{xy} = \left(\frac{17}{35}\right)(0.5\pi^2) = 2.3969$$
 36

$$k_y = \left(\frac{31}{630}\right)(0.5\pi^4) = 2.3966$$
 37

$$k_{\lambda} = \left(\frac{31}{630}\right)(0.5) = 0.02460315$$
 38

Substituting equations (35) to (38) into equation (26) gives

$$\lambda^2 = \left(\frac{2.4 + \frac{4.7938}{p^2} + \frac{2.3966}{p^4}}{\frac{31}{1260}}\right) * \frac{D}{ma^4}. That is$$

$$\lambda^2 = \left(97.5484 + \frac{194.8448}{P^2} + \frac{97.4102}{P^4}\right) * \frac{D}{ma^4} \quad 39$$

All clamped edge

rectangular

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plate

$$w_x = \sqrt{A} (R^2 - 2R^3 + R^4)$$
 40

$$w_V = \sqrt{A} (0^2 - 20^3 + 0^4)$$
 41

From equations (40) and (41), h_1 and h_2 are:

$$h_1 = R^2 - 2R^3 + R^4 42$$

$$h_2 = Q^2 - 2Q^3 + Q^4 43$$

With these we obtain:

$$\int_0^1 h_1^2 dR = \frac{1}{630} \text{ and } \int_0^1 h_2^2 dQ = \frac{1}{630}.$$

$$\int_0^1 \frac{d^4 h_1}{dR^4} h_1 dR = 0.8 \text{ and } \int_0^1 \frac{d^4 h_2}{dQ^4} h_2 dQ = 0.8$$

$$\int_0^1 \frac{d^2 h_1}{dR^2} h_1 dR = \frac{2}{105} \text{ and } \int_0^1 \frac{d^2 h_2}{dQ^2} h_2 dQ = \frac{2}{105}$$

$$k_x = (0.8) \left(\frac{1}{630}\right) = \frac{2}{1575}$$
 44

$$k_{xy} = \left(\frac{2}{105}\right) \left(\frac{2}{105}\right) = \frac{4}{11025}$$
 45

$$k_y = \left(\frac{1}{630}\right)(0.8) = \frac{2}{1575}$$
 46

$$k_{\lambda} = \left(\frac{1}{630}\right) \left(\frac{1}{630}\right) = \frac{1}{396900}$$
 47

Substituting equations (44) to (47) into equation (26) gives

$$\lambda^{2} = \left(\frac{\frac{2}{1575} + \frac{8}{11025P^{2}} + \frac{2}{1575P^{4}}}{\frac{1}{396900}}\right) * \frac{D}{ma^{4}}. That is$$

$$\lambda^{2} = \left(504 + \frac{288}{P^{2}} + \frac{504}{P^{4}}\right) * \frac{D}{ma^{4}}$$
48

RESULTS AND CONCLUSIONS

The non dimensional form of the resonating frequencies for different aspect ratios for ssss and cccc plates are shown on tables 1 and 2. A close critical examination of the tables reveals that the maximum percentage difference between the values from the present study and those from previous study is 0.09. From statistical point of view, this implies that no difference existed between the two sets of values. Thus, one can infer that the procedure, the deflection function and the energy functional formulated in this present study are reliable and adequate in CPT free vibration analysis rectangular plates. Hence, this method is recommended for stability analysis of CPT plates. It is also recommended that the present method is extended to refined plate theory analysis (RPT).

Table 1: Non dimensional form of resonating frequency of ssss isotropic thin plate

	Resonati	ng frequency,	
	$\lambda \left(\frac{1}{a^2} \sqrt{\frac{D}{m}}\right)$		
		Past	
Aspect		(Ibearugbulem	Percentage
ratio, P	Present	et al., 2014)	difference

1	19.74	19.75	-0.05
1.1	18.03	18.04	-0.06
1.2	16.73	16.73	0.00
1.3	15.71	15.72	-0.06
1.4	14.91	14.91	0.00
1.5	14.26	14.26	0.00
1.6	13.73	13.73	0.00
1.7	13.29	13.29	0.00
1.8	12.92	12.92	0.00
1.9	12.61	12.61	0.00
2	12.34	12.34	0.00

Table 2: Non dimensional form of resonating frequency of cccc isotropic thin plate

	Resonati	ng frequency,	
	$\lambda \left(\frac{1}{a^2} \sqrt{\frac{D}{m}} \right)$		
		Past	
Aspect		(Ibearugbulem	Percentage
ratio, P	Present	et al., 2014)	difference
1	36	35.97	0.08
1.1	32.96	32.93	0.09
1.2	30.77	30.75	0.07
1.3	29.17	29.15	0.07
1.4	27.97	27.95	0.07
1.5	27.05	27.03	0.07
1.6	26.33	26.31	0.08
1.7	25.77	25.75	0.08
1.8	25.32	25.30	0.08
1.9	24.95	24.94	0.04
2	24.65	24.64	0.04

REFERENCES

- Erdem, C. Imrak and Ismail Gerdemeli (2007). The problem of isotropic rectangular plate with four clamped edges. Sadhan Vol. 32, Part 3, pp. 181–186.
- Ezeh, J. C., Ibearugbulem, O. M., Njoku, K. O., and Ettu, L. O. (2013). Dynamic Analysis
 of Isotropic SSSS Plate Using Taylor Series Shape Function in Galerkin's Functional.
 International Journal of Emerging Technology and Advanced Engineering, 3 (5): 372-375.
- Hutchinson, J. R. (1992). On the bending of rectangular plates with two opposite edges simply supported. J. Appl. Mech. Trans. ASME 59: 679–681.
- Ibeabuchi, V. T. (2014). Analysis of Elastic Buckling of Stiffened Rectangular Isotropic Plates Using Virtual Work Principles. A Master's Thesis submitted to the Post graduate School, Federal University of Technology, Owerri, Nigeria.
- Ibearugbulem, O. M. (2014), Using the product of two mutually perpendicular truncated polynomial series as shape function for rectangular plate analysis, International Journal of Emerging Technologies and Engineering (IJETE) ISSN: 2348–8050, ICRTIET-2014 Conference Proceeding, 30th -31st August 2014, 1-4
- Jiu, Hui Wu, A. Q. Liu, and H. L. Chen (2007). Exact Solutions for Free-Vibration Analysis of Rectangular Plates. Journal of Applied Mechanics Vol. 74 pp. 1247-1251.
- Njoku, K. O., Ezeh, J. C., Ibearugbulem, O. M., Ettu, L. O., and Anyaogu, L. (2013). Free Vibration of Thin Rectangular Isotropic CCCC Plate Using Taylor Series Formulated Shape Function in Galerkin's Method. Academic Research International, 4 (4): 126-132.
- Szilard, R. (2004). Theories and Applications of Plate Analysis. New Jersey: John Wiley
 & Sons Inc. Taylor, R. L.and S. Govindjee (2004). Solution of clamped rectangular plate
 problems. Communi. Numer. Meth. Eng. 20: 757–765.
- Ugural, A. C. (1999). Stresses in plates and shells, 2nd ed. Singapore: McGraw-hill.
- Ventsel, E. and T. Krauthammer (2001). Thin Plates and Shells: Theory, Analysis and Applications. New York: Marcel Dekker.
- Wang, C. M., Y. C. Wang, and J. N. Reddy (2002). Problems and remedy for the Ritz method in determining stress resultant of corner supported rectangular plates. Comput.Struct. 80: 145–154
- Ye, Jianqiao (1994).Large deflection of imperfect plates by iterative BE-FE method.Journal of Engineering Mechanics, Vol. 120, No. 3 (March).